

Hubble Space Telescope observations of globular cluster systems along the Hubble sequence of spiral galaxies

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ABSTRACT

We have studied the globular cluster systems of seven giant, edge-on spiral galaxies using *Hubble Space Telescope* imaging in *V* and *I*. The galaxy sample covers the Hubble types Sa to Sc, allowing us to study the variation of the properties of globular cluster systems along the Hubble sequence. The photometry reaches ~ 1.5 mag beyond the turn-over magnitude of the globular cluster luminosity function for each galaxy. Specific frequencies of globular clusters (S_N values) were evaluated by comparing the numbers of globular clusters found in our WFPC2 pointings with those in our Milky Way that would be detected in the same spatial region if placed at the distance of the target galaxies. Results from this method were found to be consistent with the more commonly used method of constructing radial distribution functions of globular clusters. The S_N values of spirals with $B/T \lesssim 0.3$ (i.e. spirals with a Hubble type later than about Sb) are consistent with a value of $S_N = 0.55 \pm 0.25$. We suggest that this population of globular clusters represents a ‘universal’, old halo population that is present around each galaxy. Most galaxies in our sample have S_N values that are consistent with a scenario in which globular cluster systems are made up of (i) the aforementioned halo population plus (ii) a population that is associated with bulges, which grows approximately linearly with the mass of the bulge. Such scenarios include the ‘merger scenario’ for the formation of elliptical galaxies as well as the ‘multi-phase collapse’ scenario, but it seems inconsistent with the ‘secular evolution’ scenario of Pfenniger & Norman, in which bulges are formed from disc stars by means of the redistribution of angular momentum through bar instabilities and/or minor perturbations. However, there is one bulge-dominated spiral galaxy in our sample (NGC 7814) with a low S_N value that is consistent with those of the latest-type spirals. This means that the ‘secular evolution’ scenario can still be viable for some bulge-dominated spirals. Thus, our results suggest that the formation histories of galaxy bulges of early-type spirals can be significantly different from one galaxy to another.

Key words: galaxies: formation – galaxies: spiral – galaxies: star clusters.

1 INTRODUCTION

Recent observations with the *Hubble Space Telescope* (*HST*) and large-field ground-based CCD cameras have caused rapid advances in our knowledge of the formation and evolution of globular cluster

(GC) systems of galaxies. However, ellipticals and S0 galaxies have received by far the largest amount of attention, mainly because early-type galaxies usually have much richer GC systems and also suffer less from internal extinction problems than do spirals. Consequently, our knowledge of GC systems in spiral galaxies is still limited to a handful of galaxies (e.g. Harris 1991; Ashman & Zepf 1998; Kissler-Patig, Ashman & Zepf 1999). However, it should be recognized that this scarcity of data for spiral galaxy GC systems constitutes an important limitation to the use of GC systems as probes of the formation and evolution of *both* early-type *and* late-type galaxies.

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For instance, the ability to test the predictions of scenarios for the formation of galaxies and their GCs depends on our knowledge of the ‘typical’ properties of GC systems in spirals. A good example is the ‘merger model’ (Schweizer 1987; Ashman & Zepf 1992), in which elliptical galaxies are formed by mergers of spiral galaxies. In this picture, the GC systems of elliptical galaxies are composite systems. One population of GCs is associated with the progenitor spirals (i.e. metal-poor GCs with a halo-like spatial distribution), while a second population of clusters forms in the merger event (i.e. metal-rich GCs with a bulge-like spatial distribution). Ashman & Zepf (1992, 1998) described several testable predictions arising from this scenario, most of which are, however, based on comparisons with the properties of ‘typical’ spiral galaxy GC systems. A particular problem in this respect is that the characteristic specific frequency of GCs (i.e. the total number of GCs per unit galaxy luminosity) around spirals is poorly known, especially for late-type spirals (e.g. the compilation of Ashman & Zepf 1998 includes specific frequency measurements for only five spirals with Hubble type Sb or later). The absence of firm constraints on the number of GCs contributed by the progenitor spirals leads to uncertainties in many of the key predictions for elliptical galaxy GC systems that follow from the merger model. Furthermore, *metal-rich* ($-1.0 \lesssim [\text{Fe}/\text{H}] \lesssim -0.2$) GCs are known to be associated with *bulges* in spiral galaxies like the Milky Way and M31 rather than with their discs or haloes (Minniti 1995; Barmby, Holland & Huchra 2002; Forbes, Brodie & Larsen 2001), so that mergers of spirals containing significant bulges are likely to contribute metal-rich GCs to the merger remnant that are *not* formed during the merger. It is therefore very important to establish – using a significant sample of spiral galaxies – whether or not the properties of the Milky Way (MW) GC system are typical of spiral galaxy GC systems as a class, and to study the relationship between bulge luminosity and the number of metal-rich GCs in spiral galaxies.

Another important and timely area where knowledge of the properties of GC systems of spirals can yield significant progress is that of the formation and evolution of bulges of spiral galaxies. On the one hand, structural and dynamical properties of spiral bulges have long shown strong similarities with those of (low-luminosity) elliptical galaxies. Bulges follow the Fundamental Plane of elliptical galaxies (e.g. Bender, Burstein & Faber 1992), while their internal dynamics are consistent with oblate, isotropic models, just like low-luminosity ellipticals (e.g. Davies et al. 1983). Extinction-corrected optical and near-IR colours of bulges have been shown to be very similar to those of elliptical galaxies in the Coma cluster (Peletier et al. 1999), and spectroscopic metallicities and $[\text{Mg}/\text{Fe}]$ ratios in bulges are similar to those of ellipticals at a given bulge luminosity (Goudfrooij, Gorgas & Jablonka 1999). As ellipticals in rich clusters most probably formed at redshift $z > 3$ (e.g. Stanford, Eisenhardt & Dickinson 1997), these similarities argue for an early formation of bulges, in line with the original ‘monolithic collapse’ model by Eggen, Lynden-Bell & Sandage (1962). An alternative (and quite different) model is that bulges form from disc material through redistribution of angular momentum (Pfenniger & Norman 1990). In this scenario, large amounts of gas are driven into the central region of the galaxy by a stellar bar and trigger intense star formation. If enough mass is accreted, the bar itself will dissolve and the resulting galaxy will reveal a bigger bulge than before bar formation; galaxies would thus evolve from late to earlier types along the Hubble sequence (see also Pfenniger, Combes & Martinet 1994). Note that this scenario *does not involve GCs*, in that only disc *stars* would contribute to the secular building of bulges. Hence, there would be no reason for a relation between the bulge-to-total luminosity ratio

and the number of GCs per galaxy in this scenario. This provides a testable prediction, and is one of the key tests performed in the present study.

This paper is built up as follows. Section 2 describes the sample selection and the *HST* observations, while the data analysis is described in Section 3. The various results are discussed in Section 4. Finally, Section 5 summarizes the main conclusions of this study.

2 APPROACH

2.1 Galaxy sample selection

The sample consists of spiral galaxies in an edge-on configuration. We selected such galaxies in order to enable detection of GCs on both sides of the spiral discs and to minimize (spatially) the impact of dust absorption. In addition, the edge-on configuration allows one to assess radial number density distributions of the GC systems around these galaxies. Northern galaxies were selected from the Uppsala General Catalogue (UGC; Nilson 1973), while southern galaxies were selected from the Surface Photometry Catalogue of ESO-Uppsala Galaxies (ESO-LV; Lauberts & Valentijn 1989). Our two first-cut selection criteria were:

- (i) the inclination i of the galaxy ($i \geq 80^\circ$); and
- (ii) the Galactic latitude b ($b \geq 30^\circ$).

Inclinations were determined following Guthrie (1992), assuming an intrinsic flattening $q_0 \equiv (b/a)_0 = 0.11$. Inclinations were then derived using Hubble’s (1926) formula

$$\cos^2 i = (q^2 - q_0^2) / (1 - q_0^2),$$

where $q = b/a$ is the observed (catalogued) axis ratio.

In order to get the highest observing efficiency out of the *HST* observations (see below), we further considered two main (and mutually counteracting) factors. On the one hand, a galaxy should be near enough to reach $M_V = -6.0$, which is ~ 1.5 mag beyond the peak in the Galactic GC luminosity function (hereafter GCLF; Harris 1996), in a few *HST* orbits using WFPC2. On the other hand, a galaxy should be distant enough for a significant fraction of its GC system to fit in the field of view of one or two WFPC2 exposures.

From the remaining list of candidate spirals, we selected 2–3 galaxies of each main Hubble type (Sa, Sb and Sc) to study the variation of the properties of GC systems as a function of Hubble type (i.e. of bulge-to-total luminosity ratio). Global properties of the final sample galaxies are given in Table 1.

2.2 Observations

Observations were made with the Wide Field and Planetary Camera 2 (WFPC2) aboard *HST* as part of General Observer program 6685. The data consist of multiple images through the F555W and F814W filters. Our program was supplemented by archival images of a few sample galaxies (NGC 4565 and NGC 4594), taken from other *HST* programs. The images of NGC 4565 were not reanalysed; all relevant data on GC candidates in that galaxy was taken from Kissler-Patig et al. (1999). The *HST* observations are listed in Table 2 together with the exposure times for each galaxy. For a few galaxies, a subset of the images were spatially offset by $0''.5$ from the others (corresponding to an approximately integer pixel shift in both PC and WF CCDs of the WFPC2) to enable a good correction for hot pixels. The locations of the WFPC2 fields are shown in Fig. 1, superposed on to grey-scale images from the Digital Sky Survey.

Table 1. Global properties of the sample galaxies.

Galaxy	NGC 3628	NGC 4013	NGC 4517	NGC 4565	NGC 4594	IC 5176	NGC 7814	Reference(s)
RA (J2000)	11 ^h 20 ^m 16 ^s .9	11 ^h 58 ^m 31 ^s .3	12 ^h 32 ^m 45 ^s .6	12 ^h 36 ^m 20 ^s .8	12 ^h 39 ^m 59 ^s .4	22 ^h 14 ^m 55 ^s .3	00 ^h 03 ^m 14 ^s .9	1
Dec. (J2000)	+13°35′20″	+43°56′49″	+00°06′48″	+25°59′16″	−11°37′23″	−66°50′56″	+16°08′44″	1
Type (RSA)	Sb	Sb	Sc	Sb	Sa	Sbc	Sab	2
Type (T)	3.1	3.0	6.0	3.2	1.1	4.3	2.0	2
V_T	9.48	11.23	10.39	9.58	8.00	12.12 ^a	10.57	2
V_T^0	8.77	10.52	9.36	8.58	7.55	11.07 ^a	10.14	2 ^b
$(B - V)_T$	0.79	0.96	0.71	0.84	1.07	0.90 ^a	0.96	2
i	85°1	87°5	87°4	86°5	84°0	86°5	83°0	10
A_V	0.087	0.054	0.077	0.050	0.166	0.100	0.144	3 ^c
v_{hel}	843	831	1131	1229	1091	1746	1053	1
$m - M$	29.80	31.35	31.00	30.06	29.75	32.16	30.60	4,5,6,7,8,4,9 ^d
M_V^0	−21.03	−20.83	−21.64	−21.48	−22.20	−21.09	−20.46	^{b,e}
B/T	0.36	0.27	0.02	0.30	0.73	0.34	0.86	10,10,11 ^f ,12, 13,10,14 ^d
Disc r_h (arcsec)	25.8	8.4		16.6		5.4		10,10,−,12,−, 10,−
Bulge r_{dev}^g (arcsec)		5.7	9.9		50.9		32.2	−,10,11,−, 13,−,14 ^d
Bulge r_{exp} (arcsec)	22.7			20.7		3.9		10,−,−,12,−, 10,− ^d
Bulge r_{eff} (arcsec)	38.1	5.7	9.9	34.7	50.9	6.5	32.2	^h

Notes. ^aFrom Lyon Extragalactic Data base (LEDa; <http://leda.univ-lyon1.fr>). ^bCorrected for both internal and foreground reddening. ^cGalactic foreground extinction only. ^dReferences are given in galaxy order. ^eDerived from the above values. ^fCorrected value (see Section 3.2.2). ^gAll bulge radii are expressed in terms of the equivalent radius of an ellipse (see text). ^hAssumed half-light radius for the bulge (see Section 3.2.2). Equal to either r_{dev} or $1.678 r_{\text{exp}}$, depending on whether the radial light profile is fitted better with $r^{1/4}$ or with an exponential profile. References: (1) NASA/IPAC Extragalactic Data Base (NED); (2) de Vaucouleurs et al. (1991, hereafter RC3); (3) Schlegel et al. (1998); (4) Willick et al. (1997; using the Tully–Fisher relation TF); (5) Tully & Pierce (2000, using TF); (6) Ferrarese et al. (2000, using Cepheids); (7) Jacoby, Ciardullo & Harris (1996, using a variety of methods); (8) Ford et al. (1996, using Planetary Nebula Luminosity Function); (9) Tonry et al. (2001, using Surface Brightness Fluctuations); (10) this paper; (11) Gavazzi et al. (2000); (12) Wu et al. (2002); (13) Baggett et al. (1998); (14) Bothun et al. (1992).

Table 2. Summary of the *HST*/WFPC2 images of the sample galaxies.

Galaxy	<i>HST</i> program	Filter name	Exp. time (s)	N_{exp}^a	Comments
NGC 3628	6685	F555W	400	2 × 4	^{b,c}
		F814W	640	2 × 4	^{b,c}
NGC 4013	6685	F555W	800	4	^c
		F814W	1100	2	
NGC 4517	6685	F555W	400	4	^c
		F814W	640	4	^c
NGC 4565	6092	F450W	600	2 × 3	^{b,c}
		F814W	480	2 × 3	^{b,c}
NGC 4594	5512	F547M	1340	4	
		F814W	1470	6	
IC 5176	6685	F555W	800	2	
		F814W	1200	2	
NGC 7814	6685	F555W	600	2	
		F814W	600	2	

Notes. ^aNumber of exposures; ^bTwo WFPC2 pointings (exp. times listed are per pointing); ^cInteger-pixel dithered data.

2.3 Data reduction and photometry

We used the IRAF¹ package for data reduction. After standard pipeline processing and alignment of the images, we combined

the images using the IRAF/STSDAS task CRRED, thereby effectively removing both cosmic rays and hot pixels. We also trimmed each image to exclude the obscured regions near the pyramid edges of WFPC2, yielding 751×751 usable pixels per CCD. Prior to performing source photometry, the strongly varying galaxy background was fitted and subtracted. The main reason for this is to minimize errors in the photometry due to any particular choice of object aperture and sky annulus. For the CCD chips covering a significant part of the galaxies' bulges or dusty discs, we applied a median filter with a 30×30 pixel kernel to approximate the galaxy background. For the other CCDs, the smooth background gradient was fitted by a bi-cubic spline fit.

The photometry of star cluster candidates was carried out using the DAOPHOT-II (Stetson 1987) package within IRAF. The objects were selected by applying the DAOFIND task to an image prepared by dividing the data image in question by the square root of the appropriate model image (created as described in the previous paragraph). This procedure ensures uniform shot noise characteristics over the whole image. We adopted fairly tight shape constraints ($-0.6 < \text{roundness} < 0.6$; $0.2 < \text{sharpness} < 0.9$) in order to exclude extended background galaxies and faint objects distorted by noise or any residual bad pixels. The detection threshold was set at 4σ above the residual background. Although DAOFIND returned with apparent point-like detections located within – or on the edge of – dust features in the galaxy discs, we decided to exclude those from further analysis owing to the difficulty in judging the location of those sources relative to the dust features along the line of sight.

¹IRAF is distributed by the National Optical Astronomy Observatories, which is operated by the Association of Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation, USA.

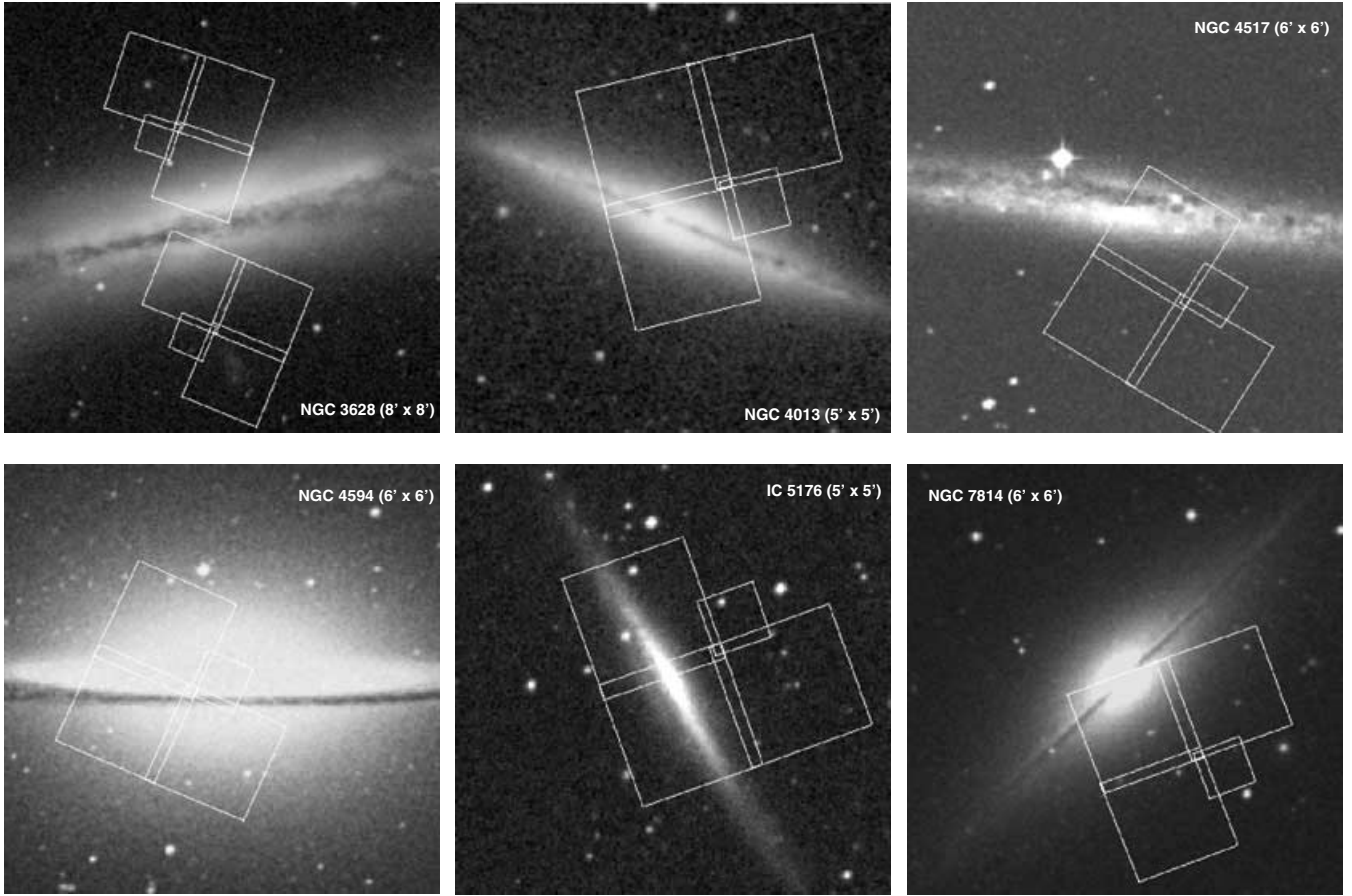


Figure 1. Images of the sample galaxies. The grey-scale images were extracted from the Digital Sky Survey, and the superposed lines show the locations of the WFPC2 field(s) of view (where each CCD chip is outlined individually). The galaxy identifications and field sizes are indicated in a corner of each image.

PSF photometry was performed using the ALLSTAR routine. For each chip, a PSF was constructed using, when possible, 10–15 bright, isolated GCs in the field. Aperture corrections were calculated using growth curve analysis on a subset of the PSF objects on each chip. One advantage of PSF photometry over the commonly-used two-pixel aperture photometry is that it avoids the systematic photometric offset due to underestimating aperture corrections for bright, resolved GCs.

One of the main sources of uncertainty in WFPC2 photometry (in general) is due to charge-transfer efficiency (CTE) problems of the CCDs for which correction recipes are available (Whitmore, Heyer & Casertano 1999; Dolphin 2000). We applied the Whitmore et al. formulae to our photometry. However, in the present case of GC photometry on a relatively high background (due to the diffuse galaxy light), the CTE correction typically affected the magnitudes by less than 0.02 mag and the colours by less than 0.01 mag – negligible for the purposes of this paper.

The measured WFPC2 STMAG magnitudes F547M, F555W and F814W were converted into Johnson–Cousins V and I magnitudes using colour terms as derived by Holtzman et al. (1995) and Goudfrooij et al. (2001). Specifically, the following conversions were used:

$$\begin{aligned} V - F547M &= (0.009 \pm 0.002) \\ &+ (0.005 \pm 0.002)(F547M - F814W) \\ &- (0.008 \pm 0.002)(F547M - F814W)^2 \end{aligned}$$

$$\begin{aligned} V - F555W &= (0.643 \pm 0.001) \\ &+ (0.004 \pm 0.002)(F555W - F814W) \\ &- (0.015 \pm 0.002)(F555W - F814W)^2 \\ I - F814W &= (-1.266 \pm 0.001) \\ &+ (0.018 \pm 0.003)(F555W - F814W) \\ &+ (0.016 \pm 0.002)(F555W - F814W)^2 \end{aligned}$$

Finally, the magnitudes of the point-like sources in the sample galaxies were corrected for Galactic foreground extinction (cf. Table 1). The A_V values were converted to A_I according to the Galactic extinction law of Rieke & Lebofsky (1985).

2.4 Cluster candidate selection

To select candidate GCs, we adopted a fairly generous selection criterion for the colour cut: $0.3 \leq V - I \leq 2.0$. The lower limit was set to 0.2 mag bluer than the bluest GC in the Milky Way, while the upper limit was set to prevent the exclusion of somewhat reddened objects (see the bottom right panel of Fig. 2). For NGC 4594, the GCs are clearly discernable on the colour–magnitude diagram (CMD; see below), so we were able to set a more stringent colour cut: $0.5 \leq V - I \leq 1.5$. We also applied a brightness cut, eliminating GC candidates brighter than G1, the most luminous M31 GC ($M_V = -10.55$, Rich et al. 1996).

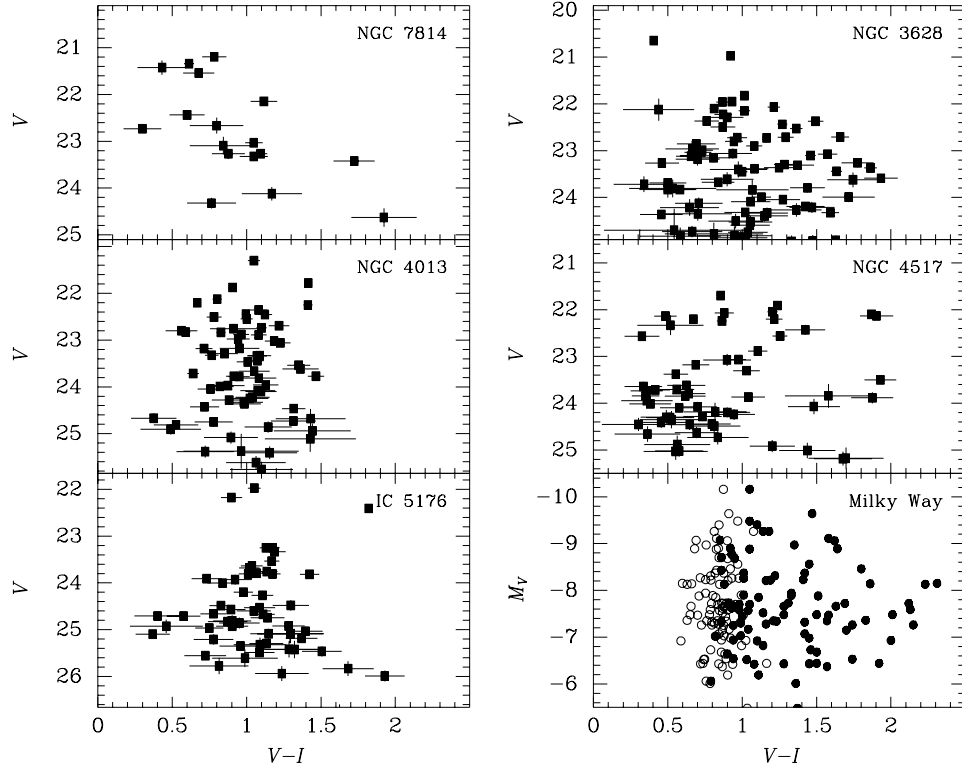


Figure 2. V versus $V - I$ colour-magnitude diagrams of candidate GCs in the five sample galaxies for which GC photometry was not published before, compared with that of the GCs in the Milky Way as taken from the 2000 update of the compilation of Harris (1996, bottom right panel). All panels cover the same range in absolute V -band magnitudes under the assumption of the galaxy distances in Table 1. Filled symbols represent observed data, whereas open symbols in the bottom right panel represent dereddened data of the Milky Way GCs. The large reddening corrections for many Milky Way GCs prompted us to use fairly generous colour cuts in selecting GC candidates for spirals with *HST* pointings encompassing the dusty discs (see Section 2.4).

To avoid complications related to excessive reddening, we excluded candidates lying within the dust lanes of their parent galaxies. Since this normally accounted for only a small fraction (~ 5 per cent) of the total detected cluster population, this should have minimal effects on our global conclusions. However, we have attempted to correct this bias in the same manner as for our incomplete spatial coverage of the target galaxies (see below).

V versus $V - I$ CMDs for the candidate GCs in five of our program galaxies² are shown in Fig. 2, together with that of the GC system of the Milky Way as a comparison. One rather unusual finding in this context is that the CMD of NGC 4517 reveals a relatively large number of GC candidates with $0.3 \leq V - I \leq 0.6$, which is bluer than Galactic GCs. However, the sizes of all these candidates (which are located on the WF chips of WFPC2) are fully consistent with GCs at the distance of NGC 4517. Hence we choose to retain those sources as GC candidates. Spectroscopy will be needed to confirm the nature of these sources.

The photometry and positions of the 50 brightest GC candidates in the five galaxies for which GC photometry was not published before are listed in Tables B1–B5 (in Appendix B).

3 RESULTS: PROPERTIES OF THE GLOBULAR CLUSTER SYSTEMS

3.1 The total number of globular clusters

The method we used to estimate the total number of GCs in the target galaxies was to make corrections for (i) contamination by Galactic foreground stars and (ii) spatial coverage. The latter was done by comparing the target GC systems to that of the Milky Way, as detailed below.

3.1.1 Correction for Galactic foreground stars

First, our estimates were corrected for foreground contamination, which was estimated from Galactic models (Bahcall & Soneira 1981). The predicted star counts at the Galactic latitude and longitude of each of the sample galaxies were subjected to the same colour and brightness cuts as the GC candidates and then subtracted from their respective bins. The total number of contaminating stars passing these criteria was small, generally ~ 3 –5 stars per pointing.

3.1.2 Correction for completeness

Artificial star experiments were performed by using the DAOPHOT-II task ADDSTAR to estimate the completeness of the finding algorithm. These were carried out in the usual fashion, using

²For the CMDs and the photometry tables of the GCs in NGC 4565 and NGC 4594, we refer to Kissler-Patig et al. (1999) and Larsen, Forbes & Brodie (2001), respectively.

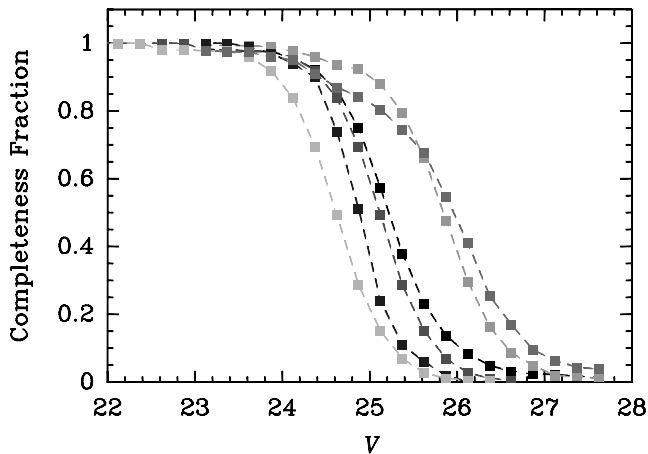


Figure 3. V-band completeness functions for the WFPC2 photometry of GC candidates. From left to right (at the 50 per cent completeness level), the curves represent the WF-chip completeness functions for NGC 7814, NGC 4594, NGC 4013, NGC 3628, NGC 4517 and IC 5176, respectively. This figure can be seen in colour in the online version of this journal on *Synergy*.

artificial GCs generated from the PSFs constructed during the photometry measurements. These ‘fake’ GCs were added in groups of 100 in randomly placed positions on each chip. For the WF chips, the 50 per cent completeness level was generally 1.0–1.5 mag beyond the turnover of the GC luminosity function (GCLF), ensuring accurate determination of the turnover point and dispersion of the GCLF. The WF-chip completeness functions are shown in Fig. 3. For each magnitude bin and each galaxy, a completeness fraction was calculated, and the number of objects in the bin was divided by this number to produce a completeness-corrected value.

3.1.3 Extrapolation over the luminosity function

In preparation for the calculation of specific GC frequencies as defined by Harris & van den Bergh (1981), the total number of GCs is defined as twice the number of GCs brighter than the turnover magnitude of the GCLF, where the GCLF is assumed to be a Gaussian (in magnitude units). We therefore fitted the completeness-corrected GCLFs by a Gaussian with a standard deviation $\sigma = 1.1$ mag (Harris 1996; Whitmore 1997) and derived the total number of clusters from the Gaussian fit.

3.1.4 Correction for incomplete spatial coverage

To estimate the total number of GCs in each galaxy, we need to correct for the incomplete spatial coverage of our data. Since each pointing only covered part of the galaxy, an extrapolation over the entire galaxy was needed to estimate the total GC population. In the case of elliptical and S0 galaxies, this extrapolation is often done by evaluating the radial distribution of GCs within the observed radial extent of the galaxy. However, the low number statistics of GCs in most spiral galaxies does not allow this to be done accurately. Hence the correction for spatial coverage was done by means of a direct comparison with the positions of GCs in the Milky Way. Since this process has been described and illustrated in detail in Kissler-Patig et al. (1999), we will only summarize the process here.

We created a mask defined by our spatial coverage of each galaxy, and applied this mask to the Milky Way. By calculating the number of objects we would detect in the Milky Way mask, if placed at the

distance of the target galaxy, we calculated the total number of GCs, N_{GC} , as follows:

$$N_{GC} = N_{OBS} \frac{N_{MW}}{N_{MSK}},$$

where N_{MSK} is the number of objects detected in the mask, and N_{MW} is the total number of GCs in the Milky Way. It should be noted that there are four different orientations of the mask that would preserve the position angle of the target galaxy (these can be seen as reflections of the mask across its horizontal or vertical axes). Our final value for N_{MSK} was the average of the values of all four possible orientations, and the standard deviation of the four values was taken into account in the uncertainty of N_{MSK} .

For this study, data regarding Milky Way GCs were taken from the McMaster catalogue (Harris 1996), containing 141 GCs with appropriate $(V - I)$ colour information. However, we adopted $N_{MW} = 160 \pm 20$ (van den Bergh 1999) as the total number of GCs in the Milky Way. The undetected Milky Way GCs are assumed to be behind the Galactic bulge. Since the scalelengths and scaleheights of the thin discs of the sample galaxies are very similar to those of the Milky Way (e.g. Bothun, Harris & Hesser 1992; Morrison, Boroson & Harding 1994; Baggett, Baggett & Anderson 1998; Gavazzi et al. 2000; Wu et al. 2002), the counterparts of such GCs would not be detected in our observations either. We therefore assume that the percentage of obscured GCs is the same in all sample galaxies.

We calculated the statistical uncertainty in the total number of GC candidates calculated this way by considering Poisson errors in the observed number of GCs, Poisson errors in the average number of Milky Way GCs in the mask (averaged over the four different orientations), and errors in the number of contaminating foreground objects. Our final numbers for the total population of GC candidates are listed in Table 3.

3.1.5 Comparison with previous studies

A potential caveat of our method to correct for incomplete spatial coverage is the implicit assumption that the spatial distribution of the GCs in the sample galaxies is similar to that of the Milky Way globular cluster system. Hence a comparison with other methods is in order. Unfortunately, there have been only very few ground-based, wide-field studies of spiral galaxy GC systems, which is probably due to the complications discussed in the Introduction. Although this is slowly changing with the use of wide-field cameras on 4-m-class telescopes (Rhode & Zepf 2002), we can currently only make a useful comparison for one galaxy: NGC 7814, for which a previous (ground-based) study of the GC system was performed by Bothun et al. (1992). Bothun et al. found a much larger number of GC candidates than we did ($N_{GC} = 498 \pm 164$ versus our 106 ± 28). From an inspection of our WFPC2 frames, we suspect that the different result is due to their counting method, which was based on object overdensities (relative to a background field) rather than colours and sizes. We find a large number of small galaxies in our WFPC2 field, which Bothun et al. (1992) probably counted along with ‘real’ GCs, while they were excluded from our GC candidate list. This once again illustrates *HST*’s unique power in discriminating GCs from small, nucleated galaxies.

A perhaps more useful comparison of our method with others can be done for the case of NGC 4594, for which Larsen et al. (2001) used the same *HST* data, but applied a different method: Larsen et al. constructed radial distribution functions of GCs to correct for incomplete spatial coverage. While that method might in principle

Table 3. Number of GCs and specific frequencies for the sample galaxies.

Galaxy (1)	GCs (detected) (2)	GCs (total) (3)	Total S_N (4)	T_{ZA93} (5)	Metal-rich GCs (all radii) (6)	Bulge S_N (all radii) (7)	Metal-rich GCs ($r < 2 r_{1/2}$) (8)	Bulge S_N ($r < 2 r_{1/2}$) (9)
NGC 3628	92 ± 10	497 ± 110	1.9 ± 0.2	3.7 ± 1.0	213 ± 48	0.83 ± 0.11	— ^a	—
NGC 4013	69 ± 8	243 ± 51	1.1 ± 0.3	2.2 ± 0.7	95 ± 20	0.44 ± 0.18	3 ± 3	0.06 ± 0.06
NGC 4517	62 ± 8	270 ± 60	0.6 ± 0.2	1.4 ± 0.5	81 ± 18	0.18 ± 0.16	— ^b	—
NGC 4565	40 ± 6	204 ± 38^c	0.6 ± 0.2	1.0 ± 0.3	122 ± 23	0.31 ± 0.15	20 ± 5^d	0.17 ± 0.15
NGC 4594	159 ± 13	1270 ± 308	1.7 ± 0.6	3.6 ± 1.1	691 ± 167	0.91 ± 0.27	429 ± 214	0.77 ± 0.44
IC 5176	57 ± 8	132 ± 25	0.5 ± 0.1	1.1 ± 0.3	67 ± 12	0.25 ± 0.09	7 ± 2^d	0.07 ± 0.07
NGC 7814	17 ± 4	106 ± 28	0.7 ± 0.2	1.5 ± 0.5	31 ± 9	0.20 ± 0.09	8 ± 5	0.06 ± 0.06

Notes. Column (1): Galaxy name. Column (2): Detected number of GC candidates in WFPC2 images of galaxy. Column (3): Total number of GC candidates around galaxy. Column (4): Total specific frequency of GC candidates. Column (5): T -parameter as defined by Zepf & Ashman (1993) as the number of GCs per unit galaxy mass (in terms of $10^9 M_\odot$). Column (6): Total number of metal-rich GCs ([Fe/H] > -1) around galaxy. Column (7): Total bulge specific frequency of metal-rich GC candidates. Column (8): Number of metal-rich GC candidates within two bulge half-light radii. Column (9): Bulge specific frequency of metal-rich GC candidates within two bulge half-light radii. ^aNot measured due to too small spatial coverage of inner bulge. ^bBulge radius too small to find bulge GCs. ^cTaken from Kissler-Patig et al. (1999). ^dNumber is twice the number of GC candidates found on the non-dusty side of the nucleus.

be expected to produce more accurate results than the Milky Way mask method used here, it requires a large population of GCs to avoid complications related to small number statistics. (In our galaxy sample, this requirement is *only* met in the case of NGC 4594.) In any case, a comparison of Larsen et al.'s results with ours is very encouraging: they find a total population of $N_{GC} = 1150 \pm 575$, compared to $N_{GC} = 1270 \pm 308$ (statistical error only) in this study. Especially heartening in this respect is that NGC 4594 (the Sombrero galaxy) is a typical Sa galaxy, which might perhaps not be expected to have a GC system particularly similar to that of the Milky Way, an Sbc galaxy.

3.2 Cluster specific frequencies

Traditionally, the specific frequency of GCs in galaxies, S_N , is defined as $S_N \equiv N_{GC} \times 10^{0.4(M_V+15)}$ (Harris & van den Bergh 1981), i.e. the total number of GCs per unit galaxy V -band luminosity (normalized to $M_V = -15$). S_N was introduced primarily for use in early-type (E and S0) galaxies, where there is little variation in stellar populations between individual galaxies. This is not necessarily the case for spirals, and indeed S_N is generally known to increase along the Hubble sequence (going from late to early-type). Averaging specific frequencies per Hubble type, the recent compilation of Ashman & Zepf (1998) supplemented by the study of two spiral galaxies by Kissler-Patig et al. (1999) yields $\langle S_N \rangle = 0.4 \pm 0.2$ for Sc spirals to 1.9 ± 0.5 for elliptical galaxies outside galaxy clusters (see Section 4). Zepf & Ashman (1993) attempted to account for differences in stellar mass-to-light ratios among galaxy Hubble types in a statistical sense by introducing a parameter T (hereafter T_{ZA93}) to be the number of GCs per unit stellar mass ($10^9 M_\odot$) of a galaxy. Conversion from luminosity to mass was achieved by assuming a characteristic M/L_V value for each galaxy Hubble type.

3.2.1 Total specific frequencies

The S_N and T_{ZA93} values of the target galaxies are listed in Table 3, calculated using the absolute magnitudes given in Table 1 and the number of GCs derived in Section 3.1. The quoted uncertainties were derived by taking into account the random errors discussed above as well as an uncertainty of 0.2 mag in the (total) absolute, dereddened V -band magnitudes of the galaxies. T_{ZA93} values were

calculated by converting the galaxy luminosities to masses following Zepf & Ashman (1993), using M/L_V values of 5.4, 6.1 and 5.0 for Sa, Sab-Sb and Sbc-Sc galaxies, respectively (cf. Faber & Gallagher 1979). We note that the uncertainty in the *distance* of a galaxy has a negligible effect on the S_N and T_{ZA93} values, since the change in the total number of GCs (due to a change in the physical area covered by the WFPC2 frames) is compensated by a similar change of the luminosity of the galaxy.

3.2.2 Bulge specific frequencies

As already mentioned in the Introduction, a view is emerging that inner metal-rich GCs in spiral galaxies may be associated with their bulges rather than with their (thick) discs. Ample evidence for this is available for the Milky Way and M 31, where the metal-rich GC systems show kinematics, metallicities, and spatial distributions matching those of the underlying bulge stars (Minniti 1995; Côté 1999; Barmby et al. 2001; Perrett et al. 2002). Such a physical association is also probably present for elliptical galaxies, where the spatial distribution of metal-rich GCs typically follows that of the spheroidal galaxy light distribution very closely, whereas the metal-poor GC system typically has a more extended distribution (e.g. Ashman & Zepf 1998, and references therein). If indeed this association can be confirmed for a large sample of spirals as well, it would provide an important causal link between the formation of a spheroidal stellar system and that of metal-rich GCs. A suggestion that this might be the case was recently provided by Forbes et al. (2001), who compared the GC systems of the Milky Way (an Sbc galaxy), M 31 (Sb) and M 104 (Sa). They argued that the ‘bulge specific frequency (bulge S_N)’ – which they defined as the number of metal-rich GCs within two bulge effective radii from the galaxy centres divided by the bulge luminosity (normalized to $M_V = -15$) – was consistent among those three galaxies (bulge $S_N \sim 0.5$), and similar to typical values for field elliptical galaxies as well.

Structural bulge properties and bulge-to-total luminosity ratios (hereafter B/T ratios) for the sample galaxies are listed in Table 1. For NGC 4517, NGC 4565, NGC 4594 and NGC 7814, these numbers were taken from the literature (respectively: Gavazzi et al. 2000; Wu et al. 2002; Baggett et al. 1998; Bothun et al. 1992). The bulge/disc decompositions for the other sample galaxies are described in Appendix A. For NGC 4013 and IC 5176, we used

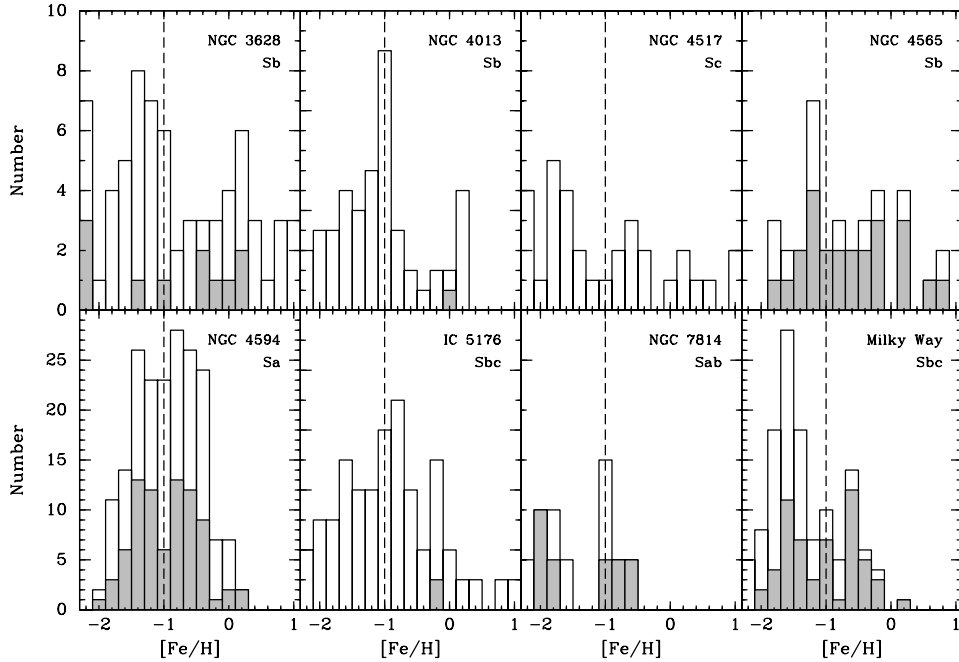


Figure 4. Metallicity distributions of GC candidates in the spiral galaxies in our sample, and of GCs in the Milky Way for comparison. The y-axis scaling is different for each galaxy, but the small tickmarks always represent 1 GC. Open histograms represent the total observed GC system, and hatched histograms represent observed GCs within 2 bulge half-light radii from the galactic centre. The vertical dashed line at $[\text{Fe}/\text{H}] = -1$ separates ‘metal-poor’ from ‘metal-rich’ GC candidates.

our own WFPC2 images (in the F555W band). For NGC 3628, we used the WIYN *V*-band image published by Howk & Savage (1999) which was graciously made available to us. Among the sample galaxies, only the bulges of NGC 4013, NGC 4517, NGC 4594 and NGC 7814 were well fitted by a de Vaucouleurs profile. The bulges of NGC 3628, NGC 4565 and IC 5176 were much better fitted by an exponential profile. Half-light radii $r_{1/2}$ for the bulges of the latter galaxies were calculated as $1.678 r_0$, where r_0 is the exponential scalelength of the profile. The final half-light bulge radii $r_{1/2}$ of all sample galaxies are listed in the last row of Table 1.

To identify ‘metal-rich’ GCs in our sample galaxies, we considered a lower limit of $[\text{Fe}/\text{H}] = -1$, which is the location of the ‘dip’ in the metallicity distribution of the Milky Way GC system (Harris 1996). This is shown in Fig. 4, superposed on to the GC metallicity histograms of the sample galaxies. The metallicities were derived by transforming the $V - I$ colours (corrected for Galactic extinction) into $[\text{Fe}/\text{H}]$ values using the recent calibration of Kissler-Patig et al. (1998). The $B - I$ colours of GCs in NGC 4565 were converted into $[\text{Fe}/\text{H}]$ values using the relation for Milky Way GCs given in Couture, Harris & Allwright (1990). Note that the $[\text{Fe}/\text{H}]$ values for GC candidates in external spiral galaxies as derived from colours are probably overestimates given the possibility of reddening by dust, especially in the inner regions. Thus, the number of ‘metal-rich’ clusters resulting from this exercise has to be regarded as an upper limit.

To allow a direct comparison with the ‘bulge S_N ’ values defined in Forbes et al. (2001), we counted the total number of ‘metal-rich’ GC candidates as well as those with a projected galactocentric radius $< 2 r_{1/2}$ identified in our images. Correction for incomplete spatial coverage was done by dividing the observed number of GCs by the fraction of the total area within $2 r_{1/2}$ sampled by our images. Since the surface density of GCs typically falls off with increasing galactocentric radius, this correction procedure was done in elliptical annuli, i.e. multiplying the number of GCs in each annulus by the

ratio of total annulus area to that which was actually observed. We chose to skip this measurement for the case of NGC 3628 since the *HST* observations of NGC 3628 only covered a very small portion of the inner bulge, rendering the extrapolated result very uncertain. In the cases of NGC 4565 and IC 5176, we only counted GCs on one side of the disc for this purpose, namely the side that was visually not significantly impacted by dust extinction. The bulge S_N values for the sample galaxies were then calculated using the B/T ratios listed in Table 1. The results are listed in Table 3.

The bulge S_N values within $2 r_{1/2}$ for most of our galaxies are roughly consistent with a value of ~ 0.1 , with the exception of NGC 4594, the Sombrero galaxy, for which we find a value of ~ 0.5 . At face value, the low bulge S_N values for most of our galaxies seem to be inconsistent with the prediction of Forbes et al. (2001), who suggested that bulge S_N values within $2 r_{1/2}$ are constant among spiral galaxies with a value around 0.5. However, systematic uncertainties of the inner bulge S_N values as derived from our data can be substantial. For example, we generally optimized our observing strategy to detect GCs in the haloes of the target galaxies (cf. Fig. 1), and dust extinction effects are strong in the innermost bulge regions. Hence we defer further discussion on the universality of the bulge S_N until high-resolution near-IR data are available. Our following discussion on trends of S_N values along the Hubble sequence will therefore focus on the ‘total’ S_N values.

4 DISCUSSION

In the context of the ‘major merger’ (Ashman & Zepf 1992) or ‘multi-phase collapse’ (Forbes, Brodie & Grillmair 1997) scenarios for the formation of early-type galaxies, the higher S_N values of ellipticals (with respect to spirals) are due to ‘extra’ GCs that were formed with high efficiency during the event that also formed the spherical stellar component (bulge). If this is indeed how bulges in spirals form in general, one would expect S_N to vary systematically

along the Hubble sequence of spirals (increasing from late-type to early-type spirals). If instead bulges were formed from disc stars by means of redistribution of angular momentum through minor perturbations and/or bar-like instabilities as in the scenario of Pfenniger & Norman (1990), one would *not* expect to see any significant variation along the Hubble sequence, since no GCs would be formed in this secular building of bulges.

Since our galaxy sample spans a large range of Hubble types (from Sa to Sc), we are in a good position to test this. We plot the S_N and T_{ZA93} values versus Hubble type and B/T ratio in Fig. 5. Values for the Milky Way and M31 (taken from Forbes et al. 2001 and references therein) are also included in that figure. There are a

number of remarkable aspects of these results, which we discuss in order below.

First of all, there does not seem to be any significant difference between trends derived from T_{ZA93} and those from S_N . Hence, we will only discuss trends involving S_N , with the understanding that the trends involving T_{ZA93} are consistent with them.

A general trend in Fig. 5 is that S_N stays basically constant at a value of ~ 0.55 for spirals with $B/T \lesssim 0.3$ (roughly corresponding to Hubble types Sb and later). As these galaxies are clearly dominated in mass by their haloes, this constant S_N (and T_{ZA93}) suggests that the formation process of massive (late-type) spiral galaxies involves the creation of a quite constant number of GCs per unit galaxy

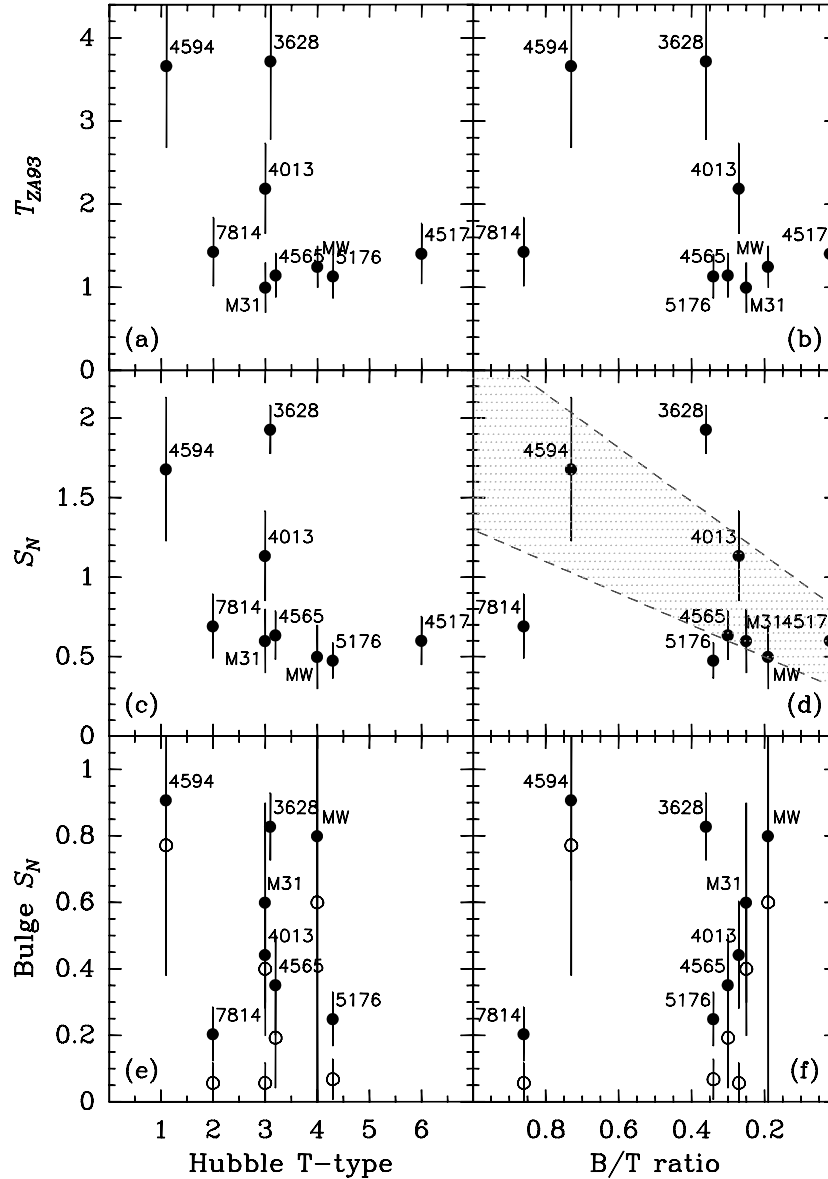


Figure 5. Relationships of the specific frequency parameters (S_N and bulge S_N) and the T parameter (introduced by Zepf & Ashman 1993; hence we identify it as T_{ZA93}) with galaxy Hubble type and bulge-to-total luminosity ratio for the sample galaxies as well as the Milky Way and M31 (data for the latter two galaxies were taken from Forbes et al. 2001). The identifications of the galaxies are plotted next to the accompanying data points (at a location least impacting the other symbols). The hatched area in panel (d) represents the region of expected S_N values for a simple model described in Section 4. In panels (e) and (f) (i.e. the bulge S_N plots), the filled circles represent the ‘total’ bulge S_N values, while the open circles represent the bulge S_N values within two half-light radii of the bulges (see Section 3.2.2). To avoid overcrowding, the galaxy identifications are only plotted along with the filled circles. This figure can be seen in colour in the online version of this journal on *Synergy*.

luminosity (and mass). We suggest that this population of GCs represents a ‘universal’, old halo population that is present around each galaxy.

It is now well established that elliptical galaxies in environments similar to that of the spirals in our sample (i.e. outside galaxy clusters) have higher specific frequencies than these late-type spirals. A χ^2 -weighted average S_N value for ellipticals outside galaxy clusters published in the literature is $S_N = 1.9 \pm 0.6$ (using S_N data from Ashman & Zepf 1998; Brown et al. 2000; Georgakakis, Forbes & Brodie 2001; Goudfrooij et al. 2001). If indeed the higher values of S_N for elliptical galaxies are due to an ‘extra’ population of GCs that was formed during the formation process of spheroidal stellar systems (such as bulges) involving a rapid dissipative collapse in which the physical conditions are such that giant star clusters are formed very efficiently (i.e. as in the ‘merger’ and ‘multi-phase collapse’ scenarios mentioned above), then one would naively expect the S_N values to scale linearly with B/T ratio from $S_N = 0.55 \pm 0.25$ at $B/T = 0$ to $S_N = 1.9 \pm 0.6$ at $B/T = 1$. This area is hatched in the middle right panel of Fig. 5.

Interestingly, the S_N values of most sample galaxies are indeed consistent with this simple prediction. There are however two galaxies for which this is clearly not the case: NGC 3628 (Sb) and NGC 7814 (Sab).

The result for NGC 3628 is quite remarkable. This Sb galaxy turns out to have an S_N value that is significantly higher than all other Sb or Sbc galaxies with S_N measurements to date. However, NGC 3628 is a peculiar galaxy. It is part of the Leo Triplet (Arp 317; Arp 1966). Large-area optical observations reveal a tidal plume extending ~ 100 kpc towards the east, as well as a bridge between NGC 3628 and NGC 3627 (Burkhead & Hutter 1981; Chromey et al. 1998), both of which contain large amounts of H I gas ($5.4 \times 10^8 M_\odot$, which is ~ 15 per cent of the total H I mass observed in the main body of NGC 3628; Haynes, Giovanelli & Roberts 1979). Rots (1978) constructed three-body orbital models for the tidal interactions between NGC 3627 and NGC 3628 that reproduce the formation of the plume and bridge, yielding an age for the plume of 8×10^8 yr since perigalacticon. This age was confirmed by photometric measurements of bright clumps in the plume by Chromey et al. (1998). It is therefore quite conceivable that the high S_N value of NGC 3628 may be due to an ‘extra’ supply of star clusters that were formed during the gaseous interaction that also caused the tidal plume and bridge. This situation seems to have occurred in M82, another nearby starburst galaxy that has undergone a past starburst triggered by tidal interactions with a gas-rich neighbour, during which a population of GCs did indeed form (see, for example, de Grijs, O’Connell & Gallagher 2001). As the gas from which such second-generation star clusters are formed is expected to be of relatively high metallicity (relative to halo GCs), this picture is consistent with the observation that the (total) bulge S_N of NGC 3628 (i.e. the number of *metal-rich* GC candidates per unit bulge luminosity) is higher than those of the other Sb–Sc galaxies in our sample by a factor similar to the total S_N . This idea should be verified by performing wide-field imaging as well as spectroscopy of a significant number of GC candidates in NGC 3628.

As to the deviant result for NGC 7814, it is perhaps somewhat surprising to find that the two earliest-type spirals in this sample (NGC 4594 and NGC 7814) have very different specific frequencies. NGC 4594 (Sa) has S_N values that are similar to those of field ellipticals (e.g. Ashman & Zepf 1998) and S0 galaxies (Kundu & Whitmore 2001), and thus seems to have formed GCs at a similarly high efficiency relative to its field stars. It is therefore highly un-

likely that bulges like the one of NGC 4594 can be formed through the ‘secular evolution’ scenario of Pfenniger & Norman (1990). On the other hand, NGC 7814 has a much lower S_N , which is consistent with the values found for the bulk of the latest-type spirals. This difference in S_N persists when considering the ‘bulge S_N ’ values. At face value, this result suggests that the (star-) formation histories of galaxy bulges of early-type spirals can be significantly different from one galaxy to another. The formation process of some bulges (such as that of NGC 4594) may have involved dissipative collapse, being accompanied by the creation of a significant number of (metal-rich) GCs, while bulges in other spiral galaxies (like NGC 7814) may have been built from the redistribution of angular momentum through bar-like instabilities as in the ‘secular evolution’ scenario. In this context, it may well be relevant that NGC 7814 is the least luminous (and thus presumably least massive) spiral galaxy in our sample, being ~ 5 times less luminous than NGC 4594. As Pfenniger (1991) showed, the accretion of small satellite galaxies (with a mass of ~ 10 per cent of that of the massive galaxy) can trigger the formation of a bar and produce a bulge through secular evolution. It is tempting to suggest that the difference in S_N value between NGC 4594 and NGC 7814 is simply due to the fact that satellite galaxies at a satellite-to-giant galaxy mass ratio of ~ 10 per cent are more common at low (giant) galaxy luminosities than at high luminosities. This suggestion is certainly consistent with observed luminosity functions of galaxies in poor groups (Jerjen, Binggeli & Freeman 2000; Flint, Bolte & Mendes de Oliveira 2001).

Given the importance of this issue in the context of our understanding of the formation of the Hubble sequence of galaxies, we argue that this suggestion should be tested with a statistically-significant sample of (edge-on) Sa–Sab spirals covering an appropriate range of luminosities. This can now be realized with only very modest amounts of observing time using the recent ACS camera aboard *HST*, which reaches some 2 mag deeper than WFPC2 at a given exposure time. The current generation of large-area CCD mosaics on ground-based 4-m-class telescopes will be very valuable as well in terms of covering whole GC populations around such galaxies in one shot.

5 SUMMARY

We have studied the GC systems of seven giant, edge-on spiral galaxies covering the Hubble types Sa to Sc, allowing us to study the variation of the properties of GC systems along the Hubble sequence. Our high-resolution *HST*/WFPC2 data (supplemented by archival WFPC2 data for two galaxies) reached ~ 1.5 mag beyond the turn-over magnitude of the GC luminosity function for each galaxy.

Specific frequencies of GCs (S_N values) were evaluated by comparing the numbers of GCs found in our WFPC2 pointings with the number of GCs in our Milky Way that would be detected in the same spatial region if placed at the distance of the target galaxies. Results from this method were compared with the more commonly used method of constructing radial distribution functions of GCs to correct for incomplete spatial coverage in the case of NGC 4594 (where both methods were possible), and found to be consistent with one another to within 1σ .

The S_N values of spirals with $B/T \lesssim 0.3$ (i.e. spirals with a Hubble type later than about Sb) are consistent with a value of $S_N = 0.55 \pm 0.25$. We suggest that this population of GCs represents a ‘universal’, old halo population that is present around each galaxy.

Most galaxies in our sample have S_N values that are consistent with a scenario in which GC systems are made up of (i) the aforementioned halo population plus (ii) a population that is associated with bulges of spirals, which grows linearly with the luminosity (and mass) of the bulge. Such scenarios include the ‘merger scenario’ for the formation of elliptical galaxies as well as the ‘multi-phase collapse’ scenario, but it seems inconsistent with the ‘secular evolution’ scenario of Pfenniger & Norman (1990), in which bulges are formed from disc stars by means of the redistribution of angular momentum through minor perturbations and/or bar instabilities.

On the other hand, the bulge-dominated spiral NGC 7814 shows a low S_N value, consistent with those of the latest-type spirals. NGC 7814 is the least luminous galaxy in our sample. Based on observed luminosity functions of galaxies in poor groups, we suggest that the ‘secular evolution’ scenario to build bulges in early-type spirals is most viable for low-luminosity spirals.

Thus, our results suggest that the formation histories of galaxy bulges of early-type spirals can be significantly different from one galaxy to another. Given the importance of our understanding of the formation of the Hubble sequence of galaxies in the context of galaxy evolution, we argue that the GC systems of a statistically-significant sample of luminous, edge-on Sa–Sb spirals be studied in the near future.

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APPENDIX A: BULGE/DISC DECOMPOSITIONS

Since bulge/disc decompositions are not the main topic of this paper, we will only provide a brief description of our method here. A more detailed description will be provided elsewhere. We first determined the scaleheights of the discs at four independent positions along their major axes, outside the region where the bulge contribution dominates. In general, the positions of the galaxy planes were determined by assuming symmetrical light distributions with respect to the disc plane, and the vertical profiles on either side of the plane were averaged together. However, the latter was not done for the bulge fits on NGC 4013 and IC 5176, since the WFPC2 frames only covered a large enough spatial extent on one side of the galaxy centre.

The vertical brightness distributions of the discs were well fitted by an exponential profile

$$\Sigma_D(z) = \Sigma_{D,0} \exp(-z/z_0),$$

where z_0 is the disc scaleheight. We found that as long as one stays outside the central areas where there is an obvious contribution from bulge light, the disc scaleheights generally do not exhibit any significant variation with distance along the major axis. This is consistent with the findings of de Grijs, Peletier & van der Kruit (1997). The average disc scaleheights are listed in Table 1.

To obtain structural parameters for the bulges of these galaxies, we used a least-squares program that fits a two-component (disc+bulge) model to the minor-axis brightness profiles. The disc component was fitted by an exponential profile and the bulge component was fitted by two different functions, namely (i) a de Vaucouleurs profile

$$\Sigma_B(z) = \Sigma_{B,e} \exp \left\{ -7.688[(z/z_e)^{1/4} - 1] \right\}, \quad (\text{A2})$$

where z_e is the effective radius along the minor axis and $\Sigma_{B,e}$ is the intensity at $z = z_e$, and (ii) a (second) exponential profile. For each galaxy, the disc scaleheight z_0 as found above was fixed in the fit. Areas where dust extinction is prominent (such as the central dust lanes) were flagged and ignored during the fitting process. The choice between the de Vaucouleurs profile and the (second) exponential profile was dictated by the χ^2 value of the fit.

Among the sample galaxies, only the bulges of NGC 4013, 4517, 4594 and 7814 were well fitted by a de Vaucouleurs profile. The bulges of the other galaxies were better fitted by an exponential profile. The effective radii (or exponential scale radii r_{exp}) of the bulges are listed in Table 1, after converting them to equivalent radii $r \equiv \sqrt{ab} = b/\sqrt{1-\epsilon}$, where a and b are the semimajor and semiminor axes of the ellipse, and ϵ its ellipticity. For NGC 3628, 4013 and 4517 and IC 5176, the bulge ellipticities were determined by using the ellipse fitting program in the ISOPHOTE package within IRAF. The dusty areas of the galaxies were flagged and ignored in the ellipse fitting process. For the galaxies with exponential bulge profiles, we also list the bulge half-light radii ($=1.678 r_{\text{exp}}$) for comparison with the effective radii of the bulges with de Vaucouleurs profiles.

APPENDIX B: TABLES OF GLOBULAR CLUSTER CANDIDATES

The photometry and astrometry of the 50 brightest GC candidates on the WFPC2 images of the galaxies for which GC

Table B1. Photometry and astrometry of the 50 brightest globular cluster candidates on the WFPC2 frames of NGC 3628. The object list is sorted by V magnitude (brightest first).

ID	RA (J2000) (^h : ^m : ^s)	Dec. (J2000) ([°] : ['] : ^{''})	V (mag)	$V - I$ (mag)
3628-1	11:20:14.64	+13:32:27.6	20.65 ± 0.01	0.41 ± 0.02
3628-2	11:20:21.12	+13:35:38.4	20.97 ± 0.02	0.92 ± 0.03
3628-3	11:20:19.92	+13:35:45.6	21.83 ± 0.02	1.01 ± 0.03
3628-4	11:20:19.44	+13:36:50.4	21.95 ± 0.01	0.93 ± 0.03
3628-5	11:20:14.40	+13:33:57.6	21.96 ± 0.03	0.87 ± 0.04
3628-6	11:20:17.76	+13:35:45.6	22.07 ± 0.03	1.21 ± 0.04
3628-7	11:20:17.76	+13:37:22.8	22.10 ± 0.02	0.81 ± 0.03
3628-8	11:20:12.72	+13:34:04.8	22.12 ± 0.24	0.44 ± 0.24
3628-9	11:20:15.60	+13:33:32.4	22.15 ± 0.02	1.01 ± 0.04
3628-10	11:20:16.56	+13:33:36.0	22.23 ± 0.02	0.87 ± 0.06
3628-11	11:20:12.72	+13:33:57.6	22.29 ± 0.08	0.90 ± 0.11
3628-12	11:20:19.92	+13:35:31.2	22.37 ± 0.05	0.76 ± 0.08
3628-13	11:20:16.80	+13:35:34.8	22.37 ± 0.05	1.49 ± 0.05
3628-14	11:20:16.32	+13:35:38.4	22.44 ± 0.03	1.27 ± 0.03
3628-15	11:20:18.00	+13:35:31.2	22.50 ± 0.07	0.87 ± 0.08
3628-16	11:20:16.80	+13:35:52.8	22.53 ± 0.04	1.36 ± 0.04
3628-17	11:20:17.04	+13:35:31.2	22.71 ± 0.06	1.66 ± 0.06
3628-18	11:20:15.84	+13:35:38.4	22.71 ± 0.04	1.29 ± 0.05
3628-19	11:20:15.84	+13:35:34.8	22.72 ± 0.04	0.97 ± 0.07
3628-20	11:20:18.48	+13:36:07.2	22.73 ± 0.02	1.16 ± 0.03
3628-21	11:20:17.76	+13:35:45.6	22.80 ± 0.04	0.94 ± 0.06
3628-22	11:20:16.08	+13:35:34.8	22.86 ± 0.10	0.69 ± 0.13
3628-23	11:20:17.28	+13:35:42.0	22.90 ± 0.04	1.08 ± 0.05
3628-24	11:20:13.44	+13:34:04.8	22.96 ± 0.17	0.66 ± 0.18
3628-25	11:20:16.08	+13:35:38.4	22.99 ± 0.05	0.73 ± 0.07
3628-26	11:20:20.40	+13:34:40.8	23.04 ± 0.05	0.70 ± 0.09
3628-27	11:20:17.28	+13:34:44.4	23.06 ± 0.11	0.94 ± 0.13
3628-28	11:20:18.24	+13:35:27.6	23.08 ± 0.06	1.57 ± 0.06
3628-29	11:20:16.80	+13:36:36.0	23.11 ± 0.02	1.46 ± 0.04
3628-30	11:20:16.32	+13:35:31.2	23.11 ± 0.06	0.68 ± 0.10
3628-31	11:20:17.28	+13:35:38.4	23.12 ± 0.07	0.66 ± 0.10
3628-32	11:20:19.20	+13:35:34.8	23.15 ± 0.09	0.81 ± 0.11
3628-33	11:20:19.92	+13:35:52.8	23.18 ± 0.14	0.70 ± 0.15
3628-34	11:20:18.00	+13:35:27.6	23.26 ± 0.10	1.78 ± 0.11
3628-35	11:20:20.40	+13:35:34.8	23.26 ± 0.08	0.46 ± 0.11
3628-36	11:20:16.56	+13:35:42.0	23.30 ± 0.06	1.28 ± 0.07
3628-37	11:20:17.04	+13:35:27.6	23.31 ± 0.08	1.37 ± 0.11
3628-38	11:20:16.56	+13:35:34.8	23.37 ± 0.08	1.25 ± 0.10
3628-39	11:20:20.40	+13:37:37.2	23.37 ± 0.03	1.86 ± 0.04
3628-40	11:20:16.32	+13:35:31.2	23.39 ± 0.06	1.08 ± 0.09
3628-41	11:20:18.96	+13:35:27.6	23.40 ± 0.17	0.97 ± 0.19
3628-42	11:20:18.72	+13:34:33.6	23.44 ± 0.03	1.63 ± 0.04
3628-43	11:20:18.48	+13:33:43.2	23.45 ± 0.03	1.00 ± 0.04
3628-44	11:20:17.04	+13:35:27.6	23.59 ± 0.10	1.93 ± 0.11
3628-45	11:20:18.48	+13:36:25.2	23.61 ± 0.15	0.90 ± 0.18
3628-46	11:20:17.04	+13:35:31.2	23.62 ± 0.16	1.75 ± 0.17
3628-47	11:20:17.52	+13:35:34.8	23.68 ± 0.08	0.84 ± 0.11
3628-48	11:20:14.88	+13:36:32.4	23.69 ± 0.06	0.50 ± 0.12
3628-49	11:20:15.84	+13:37:01.2	23.72 ± 0.17	0.34 ± 0.21
3628-50	11:20:16.80	+13:35:38.4	23.80 ± 0.10	1.44 ± 0.12

photometry has not been published before are given in Tables B1–B5. For the photometry of GC candidates in NGC 4565 and 4594, we refer to Kissler-Patig et al. (1999) and Larsen et al. (2001), respectively.

Table B2. Photometry and astrometry of the 50 brightest globular cluster candidates on the WFPC2 frames of NGC 4013. The object list is sorted by V magnitude (brightest first).

ID	RA (J2000) (^h : ^m : ^s)	Dec. (J2000) ([°] : ['] : ^{''})	V (mag)	$V - I$ (mag)
4013–1	11:58:34.32	+43:56:24.0	21.30 ± 0.01	1.05 ± 0.04
4013–2	11:58:30.96	+43:56:34.8	21.78 ± 0.02	1.41 ± 0.03
4013–3	11:58:29.52	+43:56:02.4	21.88 ± 0.02	0.91 ± 0.03
4013–4	11:58:33.36	+43:57:18.0	22.12 ± 0.02	0.80 ± 0.03
4013–5	11:58:35.28	+43:57:46.8	22.20 ± 0.01	0.67 ± 0.02
4013–6	11:58:29.76	+43:56:52.8	22.25 ± 0.02	1.41 ± 0.03
4013–7	11:58:32.64	+43:56:16.8	22.36 ± 0.04	1.08 ± 0.06
4013–8	11:58:30.24	+43:57:14.4	22.44 ± 0.02	1.00 ± 0.03
4013–9	11:58:31.44	+43:56:13.2	22.45 ± 0.02	1.12 ± 0.05
4013–10	11:58:33.84	+43:57:25.2	22.51 ± 0.03	0.78 ± 0.05
4013–11	11:58:28.56	+43:56:52.8	22.55 ± 0.02	1.00 ± 0.04
4013–12	11:58:31.68	+43:57:00.0	22.69 ± 0.05	1.22 ± 0.07
4013–13	11:58:33.36	+43:56:45.6	22.74 ± 0.02	1.10 ± 0.03
4013–14	11:58:34.08	+43:56:24.0	22.76 ± 0.10	0.91 ± 0.16
4013–15	11:58:25.44	+43:57:43.2	22.80 ± 0.10	0.56 ± 0.11
4013–16	11:58:31.44	+43:58:08.4	22.83 ± 0.02	0.59 ± 0.03
4013–17	11:58:30.00	+43:57:21.6	22.83 ± 0.01	0.83 ± 0.03
4013–18	11:58:29.52	+43:56:56.4	22.88 ± 0.04	0.96 ± 0.05
4013–19	11:58:32.40	+43:56:34.8	22.89 ± 0.04	1.08 ± 0.05
4013–20	11:58:27.36	+43:56:09.6	22.98 ± 0.03	0.94 ± 0.08
4013–21	11:58:32.88	+43:56:27.6	23.02 ± 0.04	1.18 ± 0.06
4013–22	11:58:30.96	+43:57:00.0	23.06 ± 0.06	1.23 ± 0.07
4013–23	11:58:33.36	+43:55:51.6	23.18 ± 0.10	0.95 ± 0.13
4013–24	11:58:31.68	+43:57:10.8	23.18 ± 0.04	0.71 ± 0.06
4013–25	11:58:29.76	+43:57:10.8	23.29 ± 0.08	0.85 ± 0.09
4013–26	11:58:31.20	+43:57:57.6	23.32 ± 0.03	0.77 ± 0.05
4013–27	11:58:32.88	+43:56:34.8	23.33 ± 0.05	1.09 ± 0.08
4013–28	11:58:28.80	+43:56:13.2	23.34 ± 0.05	1.07 ± 0.07
4013–29	11:58:32.40	+43:57:57.6	23.44 ± 0.04	1.07 ± 0.04
4013–30	11:58:31.20	+43:57:03.6	23.46 ± 0.06	1.01 ± 0.08
4013–31	11:58:32.64	+43:56:42.0	23.54 ± 0.06	1.35 ± 0.08
4013–32	11:58:31.92	+43:56:42.0	23.62 ± 0.06	1.36 ± 0.12
4013–33	11:58:29.52	+43:56:49.2	23.66 ± 0.06	1.05 ± 0.10
4013–34	11:58:24.72	+43:57:25.2	23.71 ± 0.03	0.64 ± 0.04
4013–35	11:58:30.72	+43:57:39.6	23.77 ± 0.02	1.47 ± 0.06
4013–36	11:58:33.84	+43:57:32.4	23.77 ± 0.04	0.91 ± 0.05
4013–37	11:58:31.20	+43:56:16.8	23.78 ± 0.04	0.95 ± 0.07
4013–38	11:58:30.72	+43:56:34.8	23.81 ± 0.07	1.08 ± 0.12
4013–39	11:58:33.84	+43:56:20.4	23.96 ± 0.05	1.13 ± 0.09
4013–40	11:58:28.56	+43:57:03.6	23.97 ± 0.08	0.87 ± 0.11
4013–41	11:58:30.00	+43:57:07.2	23.99 ± 0.05	0.82 ± 0.09
4013–42	11:58:32.40	+43:57:14.4	24.04 ± 0.06	0.76 ± 0.09
4013–43	11:58:30.24	+43:56:56.4	24.08 ± 0.09	1.08 ± 0.12
4013–44	11:58:33.12	+43:56:38.4	24.10 ± 0.08	1.10 ± 0.10
4013–45	11:58:31.68	+43:58:15.6	24.22 ± 0.10	1.04 ± 0.11
4013–46	11:58:28.80	+43:56:52.8	24.25 ± 0.08	1.02 ± 0.12
4013–47	11:58:34.08	+43:56:42.0	24.28 ± 0.09	0.88 ± 0.12
4013–48	11:58:33.84	+43:56:45.6	24.33 ± 0.07	0.98 ± 0.13
4013–49	11:58:31.92	+43:56:34.8	24.36 ± 0.06	0.99 ± 0.11
4013–50	11:58:33.84	+43:57:28.8	24.43 ± 0.05	0.72 ± 0.10

Table B3. Photometry and astrometry of the 50 brightest globular cluster candidates on the WFPC2 frames of NGC 4517. The object list is sorted by V magnitude (brightest first).

ID	RA (J2000) (^h : ^m : ^s)	Dec. (J2000) ([°] : ['] : ^{''})	V (mag)	$V - I$ (mag)
4517–1	12:32:43.68	+00:06:07.2	21.70 ± 0.02	0.86 ± 0.02
4517–2	12:32:47.52	+00:06:00.0	21.91 ± 0.02	1.24 ± 0.02
4517–3	12:32:43.68	+00:06:32.4	22.05 ± 0.03	1.20 ± 0.03
4517–4	12:32:42.48	+00:05:16.8	22.07 ± 0.05	0.88 ± 0.06
4517–5	12:32:48.72	+00:05:31.2	22.10 ± 0.02	1.87 ± 0.03
4517–6	12:32:46.32	+00:05:16.8	22.13 ± 0.02	1.90 ± 0.11
4517–7	12:32:42.96	+00:06:03.6	22.14 ± 0.07	0.48 ± 0.08
4517–8	12:32:43.44	+00:07:08.4	22.20 ± 0.04	1.22 ± 0.06
4517–9	12:32:42.72	+00:05:38.4	22.21 ± 0.02	0.67 ± 0.02
4517–10	12:32:46.32	+00:06:28.8	22.24 ± 0.03	0.86 ± 0.03
4517–11	12:32:41.52	+00:06:25.2	22.33 ± 0.21	0.52 ± 0.22
4517–12	12:32:44.16	+00:06:32.4	22.43 ± 0.11	1.42 ± 0.13
4517–13	12:32:43.92	+00:06:32.4	22.56 ± 0.04	1.26 ± 0.05
4517–14	12:32:43.20	+00:06:28.8	22.57 ± 0.09	0.33 ± 0.12
4517–15	12:32:42.96	+00:06:21.6	22.89 ± 0.06	1.11 ± 0.06
4517–16	12:32:43.68	+00:06:21.6	23.07 ± 0.03	0.98 ± 0.05
4517–17	12:32:46.08	+00:06:28.8	23.08 ± 0.13	0.90 ± 0.16
4517–18	12:32:44.40	+00:06:25.2	23.18 ± 0.05	0.69 ± 0.09
4517–19	12:32:41.04	+00:04:58.8	23.31 ± 0.06	1.03 ± 0.09
4517–20	12:32:48.48	+00:05:27.6	23.38 ± 0.07	0.55 ± 0.08
4517–21	12:32:47.76	+00:05:24.0	23.50 ± 0.06	1.93 ± 0.11
4517–22	12:32:45.12	+00:06:25.2	23.62 ± 0.10	0.63 ± 0.13
4517–23	12:32:46.32	+00:06:36.0	23.64 ± 0.10	0.34 ± 0.13
4517–24	12:32:45.36	+00:06:36.0	23.71 ± 0.08	0.56 ± 0.12
4517–25	12:32:44.88	+00:06:25.2	23.72 ± 0.07	0.42 ± 0.10
4517–26	12:32:43.20	+00:07:08.4	23.73 ± 0.09	0.36 ± 0.11
4517–27	12:32:45.84	+00:06:21.6	23.78 ± 0.12	0.74 ± 0.15
4517–28	12:32:45.36	+00:06:32.4	23.80 ± 0.09	0.64 ± 0.11
4517–29	12:32:43.20	+00:07:12.0	23.84 ± 0.07	0.35 ± 0.09
4517–30	12:32:45.60	+00:04:55.2	23.84 ± 0.25	1.58 ± 0.29
4517–31	12:32:45.84	+00:06:28.8	23.85 ± 0.10	0.62 ± 0.14
4517–32	12:32:45.36	+00:06:32.4	23.86 ± 0.08	0.32 ± 0.11
4517–33	12:32:45.60	+00:06:25.2	23.87 ± 0.08	1.04 ± 0.11
4517–34	12:32:43.44	+00:06:21.6	23.89 ± 0.12	1.88 ± 0.14
4517–35	12:32:43.20	+00:06:32.4	23.92 ± 0.07	1.69 ± 0.08
4517–36	12:32:45.60	+00:06:25.2	23.95 ± 0.13	0.36 ± 0.17
4517–37	12:32:44.40	+00:06:25.2	24.02 ± 0.08	0.38 ± 0.14
4517–38	12:32:42.00	+00:06:25.2	24.07 ± 0.16	1.48 ± 0.18
4517–39	12:32:45.84	+00:06:36.0	24.08 ± 0.11	0.70 ± 0.13
4517–40	12:32:44.40	+00:06:21.6	24.10 ± 0.10	0.58 ± 0.13
4517–41	12:32:42.96	+00:06:14.4	24.19 ± 0.19	0.82 ± 0.21
4517–42	12:32:44.40	+00:06:10.8	24.20 ± 0.11	0.90 ± 0.15
4517–43	12:32:43.44	+00:06:25.2	24.23 ± 0.19	0.65 ± 0.22
4517–44	12:32:45.84	+00:06:36.0	24.24 ± 0.11	0.95 ± 0.14
4517–45	12:32:45.84	+00:06:32.4	24.28 ± 0.08	0.73 ± 0.13
4517–46	12:32:45.84	+00:06:36.0	24.30 ± 0.16	0.49 ± 0.17
4517–47	12:32:43.20	+00:06:32.4	24.31 ± 0.11	0.52 ± 0.16
4517–48	12:32:44.40	+00:06:25.2	24.35 ± 0.19	0.53 ± 0.21
4517–49	12:32:45.60	+00:06:32.4	24.41 ± 0.12	0.46 ± 0.18
4517–50	12:32:45.36	+00:06:28.8	24.45 ± 0.11	0.80 ± 0.13

Table B4. Photometry and astrometry of the 50 brightest globular cluster candidates on the WFPC2 frames of IC 5176. The object list is sorted by V magnitude (brightest first).

ID	RA (J2000) (^h : ^m : ^s)	Dec. (J2000) ([°] : ['] : ^{''})	V (mag)	$V - I$ (mag)
5176-1	22:14:54.96	-66:50:38.4	21.97 ± 0.04	1.05 ± 0.04
5176-2	22:14:43.44	-66:51:03.6	22.18 ± 0.06	0.90 ± 0.07
5176-3	22:15:00.00	-66:51:14.4	22.41 ± 0.02	1.82 ± 0.02
5176-4	22:15:05.04	-66:50:09.6	23.25 ± 0.04	1.18 ± 0.05
5176-5	22:14:51.84	-66:51:18.0	23.25 ± 0.04	1.13 ± 0.04
5176-6	22:15:01.68	-66:50:34.8	23.33 ± 0.06	1.19 ± 0.08
5176-7	22:14:56.88	-66:51:25.2	23.54 ± 0.05	1.17 ± 0.05
5176-8	22:15:00.96	-66:50:42.0	23.63 ± 0.04	1.04 ± 0.07
5176-9	22:14:52.08	-66:51:00.0	23.68 ± 0.04	1.02 ± 0.07
5176-10	22:14:51.12	-66:51:32.4	23.76 ± 0.06	1.14 ± 0.07
5176-11	22:14:50.16	-66:51:21.6	23.79 ± 0.04	1.07 ± 0.06
5176-12	22:14:51.36	-66:51:21.6	23.81 ± 0.05	1.01 ± 0.06
5176-13	22:14:57.12	-66:51:18.0	23.81 ± 0.04	1.17 ± 0.05
5176-14	22:14:54.96	-66:51:32.4	23.82 ± 0.05	1.42 ± 0.06
5176-15	22:15:02.88	-66:49:51.6	23.91 ± 0.08	0.73 ± 0.10
5176-16	22:14:58.08	-66:51:39.6	23.93 ± 0.11	0.92 ± 0.11
5176-17	22:14:59.76	-66:49:55.2	24.00 ± 0.09	0.84 ± 0.10
5176-18	22:15:02.64	-66:50:09.6	24.20 ± 0.06	0.98 ± 0.07
5176-19	22:14:52.80	-66:51:00.0	24.27 ± 0.06	1.11 ± 0.07
5176-20	22:14:53.52	-66:50:56.4	24.48 ± 0.11	1.30 ± 0.12
5176-21	22:14:58.56	-66:49:55.2	24.49 ± 0.06	0.83 ± 0.08
5176-22	22:14:56.64	-66:51:18.0	24.53 ± 0.09	1.09 ± 0.10
5176-23	22:14:51.36	-66:51:18.0	24.57 ± 0.07	0.89 ± 0.10
5176-24	22:14:51.60	-66:51:07.2	24.60 ± 0.09	1.05 ± 0.11
5176-25	22:14:56.64	-66:51:39.6	24.66 ± 0.08	0.78 ± 0.09
5176-26	22:15:03.84	-66:50:20.4	24.67 ± 0.06	1.11 ± 0.08
5176-27	22:14:53.52	-66:50:56.4	24.71 ± 0.10	0.40 ± 0.16
5176-28	22:14:51.84	-66:51:18.0	24.71 ± 0.09	0.58 ± 0.11
5176-29	22:15:00.24	-66:50:38.4	24.74 ± 0.10	1.14 ± 0.11
5176-30	22:14:56.40	-66:51:25.2	24.81 ± 0.12	0.90 ± 0.13
5176-31	22:14:54.72	-66:50:42.0	24.83 ± 0.11	0.87 ± 0.13
5176-32	22:14:56.64	-66:50:24.0	24.86 ± 0.12	0.95 ± 0.14
5176-33	22:15:01.92	-66:50:13.2	24.92 ± 0.07	1.28 ± 0.11
5176-34	22:14:55.92	-66:50:24.0	24.92 ± 0.09	0.90 ± 0.11
5176-35	22:14:53.04	-66:50:52.8	24.93 ± 0.14	0.46 ± 0.24
5176-36	22:14:51.12	-66:51:18.0	24.97 ± 0.08	0.75 ± 0.09
5176-37	22:15:05.04	-66:50:09.6	25.03 ± 0.10	1.40 ± 0.11
5176-38	22:14:55.68	-66:50:31.2	25.09 ± 0.20	1.30 ± 0.22
5176-39	22:14:55.68	-66:50:31.2	25.09 ± 0.20	1.29 ± 0.22
5176-40	22:14:51.84	-66:51:21.6	25.09 ± 0.11	1.15 ± 0.12
5176-41	22:15:04.08	-66:49:48.0	25.10 ± 0.09	0.37 ± 0.12
5176-42	22:14:54.96	-66:51:36.0	25.18 ± 0.09	1.37 ± 0.11
5176-43	22:14:50.64	-66:51:18.0	25.21 ± 0.11	0.78 ± 0.13
5176-44	22:15:01.68	-66:50:45.6	25.30 ± 0.09	1.13 ± 0.11
5176-45	22:15:01.20	-66:50:45.6	25.32 ± 0.15	1.09 ± 0.18
5176-46	22:14:59.52	-66:49:58.8	25.35 ± 0.10	0.96 ± 0.14
5176-47	22:15:01.68	-66:50:56.4	25.42 ± 0.17	1.28 ± 0.19
5176-48	22:14:52.56	-66:50:13.2	25.42 ± 0.19	1.32 ± 0.19
5176-49	22:14:51.36	-66:51:21.6	25.46 ± 0.11	1.51 ± 0.13
5176-50	22:14:54.72	-66:50:20.4	25.49 ± 0.09	1.09 ± 0.10

Table B5. Photometry and astrometry of the globular cluster candidates on the WFPC2 frames of NGC 7814. The object list is sorted by V magnitude (brightest first).

ID	RA (J2000) (^h : ^m : ^s)	Dec. (J2000) ([°] : ['] : ^{''})	V (mag)	$V - I$ (mag)
7814-1	00:03:13.20	+16:08:27.6	21.19 ± 0.03	0.78 ± 0.08
7814-2	00:03:09.84	+16:06:50.4	21.34 ± 0.03	0.61 ± 0.03
7814-3	00:03:10.80	+16:08:02.4	21.43 ± 0.15	0.43 ± 0.17
7814-4	00:03:08.64	+16:08:49.2	21.54 ± 0.10	0.68 ± 0.10
7814-5	00:03:13.44	+16:08:24.0	22.15 ± 0.08	1.12 ± 0.09
7814-6	00:03:12.72	+16:08:20.4	22.44 ± 0.07	0.60 ± 0.12
7814-7	00:03:11.28	+16:07:55.2	22.67 ± 0.17	0.80 ± 0.18
7814-8	00:03:12.24	+16:08:42.0	22.73 ± 0.06	0.30 ± 0.13
7814-9	00:03:12.48	+16:07:26.4	23.03 ± 0.04	1.05 ± 0.06
7814-10	00:03:10.80	+16:08:45.6	23.09 ± 0.19	0.84 ± 0.22
7814-11	00:03:14.16	+16:06:50.4	23.27 ± 0.10	0.88 ± 0.11
7814-12	00:03:13.20	+16:06:57.6	23.27 ± 0.03	1.09 ± 0.04
7814-13	00:03:15.12	+16:08:13.2	23.32 ± 0.08	1.05 ± 0.09
7814-14	00:03:12.96	+16:07:30.0	23.42 ± 0.03	1.72 ± 0.14
7814-15	00:03:13.68	+16:08:24.0	24.12 ± 0.14	1.17 ± 0.20
7814-16	00:03:13.44	+16:08:24.0	24.32 ± 0.11	0.76 ± 0.16
7814-17	00:03:13.44	+16:08:34.8	24.63 ± 0.20	1.92 ± 0.22

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